ELEN E3401: Electromagnetics

Spring 2025

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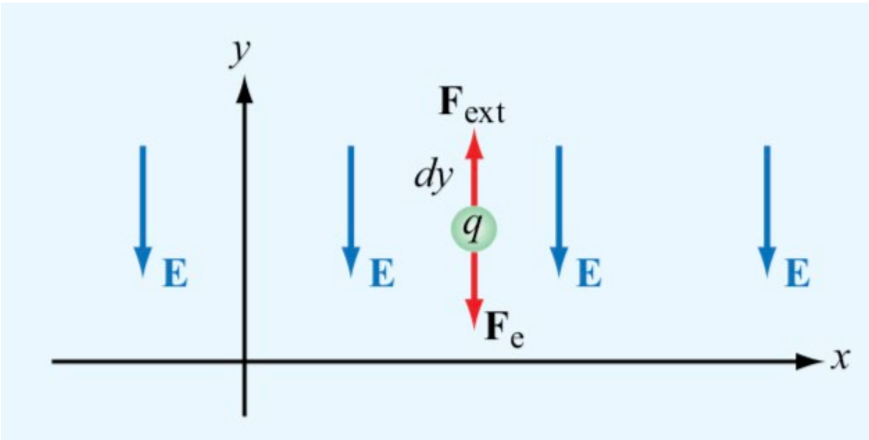
Lecture #11

COLUMBIA | ENGINEERING
The Fu Foundation School of Engineering and Applied Science

Electric Potential (voltage, voltage potential)



Consider a positive charge, q in uniform $\vec{E}$ field:
$\vec{E} = -\hat{y}E$ in the –y direction

Electric field will exert force:
$\vec{F}_e = q\vec{E}$ in the -y direction

To move the charge in the +y direction, at constant speed (no acceleration) – need to exert force, $\vec{F}_{ext} = -\vec{F}_e = -q\vec{E} \rightarrow \vec{F}_{ext} + \vec{F}_e = 0$

Work done, or energy expended to move any object a distance, *dl*

$$dW = \vec{F} \cdot d\vec{l}$$

Work done on charge: (in Joules)

$$\underbrace{dW}_{\text{Differential electric potential energy}} = \vec{F}_{ext} \cdot d\vec{l} = -q\vec{E} \cdot d\vec{l} = -q(-\hat{y}E) \cdot \hat{y}dy = qEdy$$

Electric Potential (voltage, voltage potential)

Differential electric potential energy / unit charge = dV

$$dV = \frac{dW}{q} = -\vec{E} \cdot d\vec{l} \;\; (\frac{J}{C} \; or \; V) \qquad 1V = 1\frac{J}{C} \;\; then \; \vec{E}: \frac{V}{m}$$

Potential difference: moving point charge from any point P_1 to P_2 along any path:

$$\int_{P_1}^{P_2} dV = -\int_{P_1}^{P_2} \vec{E} \cdot d\vec{l}$$

$$V_{21} = V_2 - V_1 = -\int_{P_1}^{P_2} \vec{E} \cdot d\vec{l}$$

V_1, V_2 : electric potentials at P_1 , P_2 **V_{21} : independent of specific path**

Electric Potential

KVL: net voltage drop around closed loop is zero

From above: $-\int \vec{E} \cdot d\vec{l}$

→ if we go from P_1 to P_2 then return P_2 to P_1, net voltage = 0

$$\oint_C \vec{E} \cdot d\vec{l} = 0 \qquad \text{Electrostatics}$$

Line integral of electrostatic field, $\vec{E}$ around any closed contour, C is zero

→ Conservative vector field = irrotational

Electric Potential

→ If $\vec{E}$ is <u>time-varying</u> → no longer conservative

If $\frac{\partial}{\partial t} = 0$ → $\vec{\nabla} \times \vec{E} = 0$ $\qquad$ $\vec{\nabla} \times \vec{E} = -\frac{\partial \vec{B}}{\partial t}$

Surface integral of $(\vec{\nabla} \times \vec{E})$ over open surface, S, then apply Stoke's Theorem

$$\int_S (\vec{\nabla} \times \vec{E}) \cdot d\vec{S} = \oint_C \vec{E} \cdot d\vec{l} = 0$$

<u>conservative</u>

If $\frac{\partial}{\partial t} = 0$

Electric Potential

We take reference at ∞ (similar to ground) → $V_1 = 0$

At point, P:

$$V = -\int_{\infty}^{P} \vec{E} \cdot d\vec{l}$$

Electric potential of point charge: $\vec{E} = \hat{R}\dfrac{q}{4\pi\epsilon R^2}$ Choose $d\vec{l} = \hat{R}dR$

$$V = -\int_{\infty}^{R} \hat{R}\frac{q}{4\pi\epsilon R^2} \cdot \hat{R}dR = \frac{q}{4\pi\epsilon R}\ [V]$$

If charge located at position vector, $\vec{R}_1$: $V = \dfrac{q}{4\pi\epsilon|\vec{R} - \vec{R}_1|}\ [V]$

Collection of charges: $V = \dfrac{1}{4\pi\epsilon}\displaystyle\sum_{i=1}^{N}\dfrac{q_i}{|\vec{R} - \vec{R}_i|}\ [V]$

Electric Potential

We can obtain the electric potential for charge distributions:

$$V = \frac{1}{4\pi\epsilon}\int_{\mathcal{V}'} \frac{\rho_V}{R'}\, d\mathcal{V}'$$ (Volume distribution)

$$V = \frac{1}{4\pi\epsilon}\int_{S'} \frac{\rho_s}{R'}\, ds'$$ (Surface distribution)

$$V = \frac{1}{4\pi\epsilon}\int_{l'} \frac{\rho_l}{R'}\, dl'$$ (Line distribution)

Electric Field → given electric potential

$$dV = -\vec{E} \cdot d\vec{l}$$

For any scalar function $dT = \vec{\nabla} T \cdot d\vec{l}$

gradient

Express $dV = \vec{\nabla} V \cdot d\vec{l}$

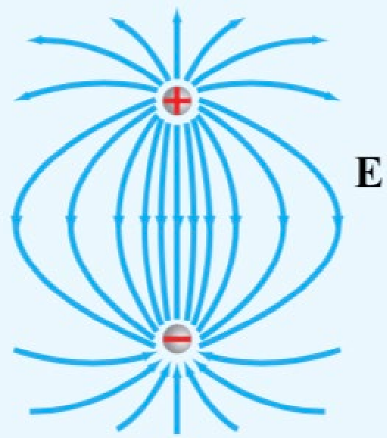
$$\vec{E} = -\vec{\nabla} V$$

We can now determine $\vec{E}$ by obtaining V for <u>any charge distribution</u>
→ then obtaining $-\vec{\nabla} V$ to get $\vec{E}$

This approach is computationally better than direct integral of Coulomb's law to obtain $\vec{E}$

Electric dipole

Model: 2 point charges of opposite polarity



(b) Electric-field pattern

(will show in PS #4)

Dipole moment: $\vec{p} = q\vec{d}$ $\vec{d}$ points from $-q$ to $+q$

Electric dipole potential: $V = \dfrac{\vec{p} \cdot \hat{R}}{4\pi\epsilon_0 R^2}$

$$\vec{E} = -\vec{\nabla}V = \frac{qd}{4\pi\epsilon_0 R^3}\left(\hat{R}2\cos(\theta) + \hat{\theta}\sin(\theta)\right) \left[\frac{V}{m}\right]$$

Poisson's Equation

$\vec{D} = \epsilon\,\vec{E}$ Differential Gauss's Law: $\vec{\nabla} \cdot \vec{E} = \frac{\rho_V}{\epsilon}$

$\vec{E} = -\vec{\nabla} V$ Insert → $\vec{\nabla} \cdot \left(\vec{\nabla} V\right) = \frac{-\rho_V}{\epsilon}$

$$\vec{\nabla} \cdot \left(\vec{\nabla} V\right) = \nabla^2 V = \frac{\partial^2 V}{\partial x^2} + \frac{\partial^2 V}{\partial y^2} + \frac{\partial^2 V}{\partial z^2}$$ Laplacian of scalar function

$$\nabla^2 V = \frac{-\rho_V}{\epsilon}$$ ← Poisson's equation

Potential $V = \dfrac{1}{4\pi\epsilon} \displaystyle\int_{\mathcal{V}'} \frac{\rho_V}{R'}\, d\mathcal{V}'$

If no charges inside: $\nabla^2 V = 0$ Laplace's equation

Maxwell's Equations: statics

<u>Electrostatics</u>:

$$\vec{\nabla} \cdot \vec{D} = \rho_V$$

$$\vec{\nabla} \times \vec{E} = 0$$

<u>Magnetostatics:</u>

$$\vec{\nabla} \cdot \vec{B} = 0$$

$$\vec{\nabla} \times \vec{H} = \vec{J}$$

Conductors

Free electrons loosely attached to atoms.
With applied $\vec{E}$ external field, electrons migrate in opposite to field direction

Conduction current: $\vec{J} = \sigma\vec{E}$ (A/m^2) Ohm's Law

Silver:	6.2×10^7	[S/m] [σ increases with decreasing temperature]
Gold:	4.1×10^7	
Germanium:	2.2	Semiconductors → need to dope
Silicon:	4.4×10^{-4}	
Glass:	10^{-12}	→ dielectric

Conductors

Perfect dielectric $\sigma = 0$ Perfect conductor $\sigma = \infty$

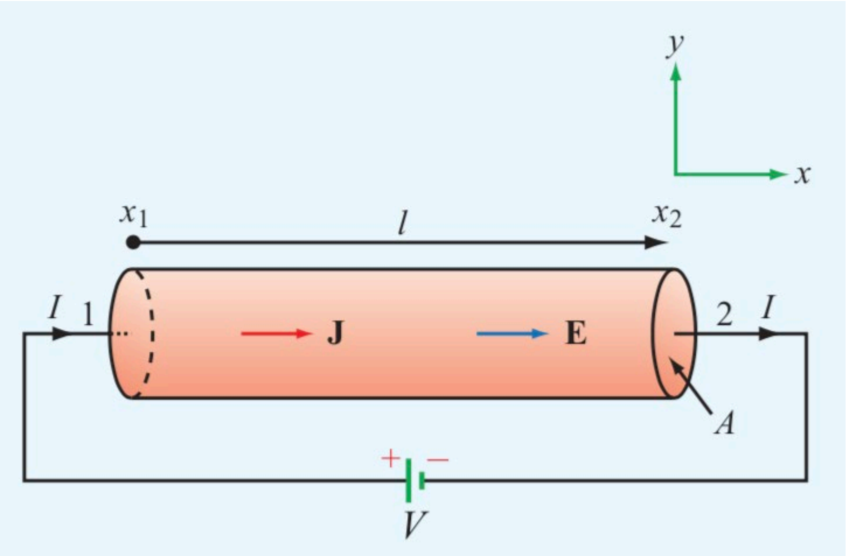
$\vec{J} = \sigma\vec{E}$ (A/m^2) Ohm's Law

$$\begin{cases} \text{With } \sigma = 0 \text{ (dielectric), } \vec{J} = 0 \text{ (independent of } \vec{E}\text{)} \\ \text{With } \sigma = \infty,\ \vec{E} = \dfrac{\vec{J}}{\sigma} = 0 \end{cases}$$

Perfect dielectric $\longrightarrow$ $\vec{J} = 0$ Perfect conductor (Equipotential) $\longrightarrow$ $\vec{E} = 0$

$$V_{21} = -\int_{P_1}^{P_2} \vec{E} \cdot d\vec{l} = V_2 - V_1 = 0$$

Resistance



Conductor length: $l = x_2 - x_1$
Cross section area: A

$$\vec{E} = \hat{x}E_x$$

Higher potential point 1
Lower potential point 2

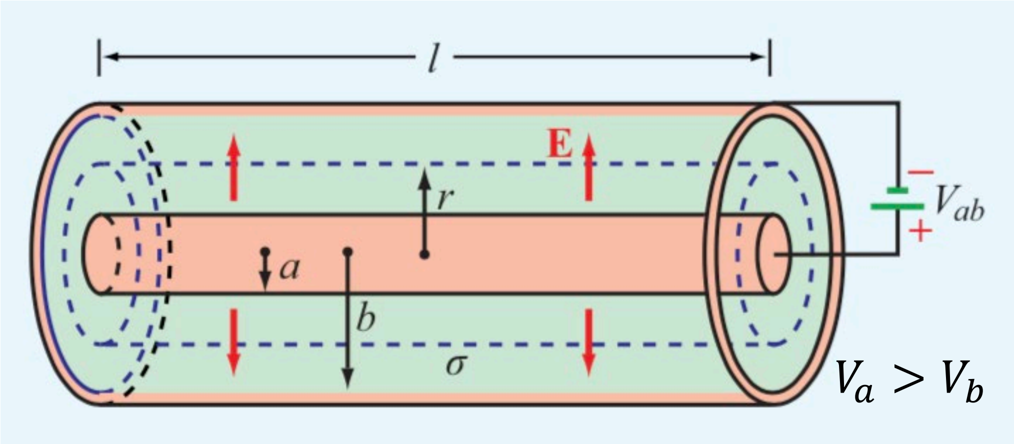
$$V = V_1 - V_2 = -\int_{x_2}^{x_1} \vec{E} \cdot d\vec{l} = -\int_{x_2}^{x_1} \hat{x}E_x \cdot \hat{x}dl = E_x l$$

$$I = \int_A \vec{J} \cdot d\vec{s} = \int_A \sigma\vec{E} \cdot d\vec{s} = \sigma E_x A \qquad R = \frac{V}{I} = \frac{E_x l}{\sigma E_x A} = \frac{l}{\sigma A}\ [\Omega]$$

Can generalize result for resistance

$$R = \frac{V}{I} = \frac{-\int_l \vec{E} \cdot d\vec{l}}{\int_S \vec{J} \cdot d\vec{s}} = \frac{-\int_l \vec{E} \cdot d\vec{l}}{\int_S \sigma\vec{E} \cdot d\vec{s}} \qquad R = \frac{l}{\sigma A}\ [\Omega] \qquad G = \frac{1}{R} = \frac{\sigma A}{l}\ [S]$$

Example - coax cable



Conductance of coax

Voltage inner conductor > outer

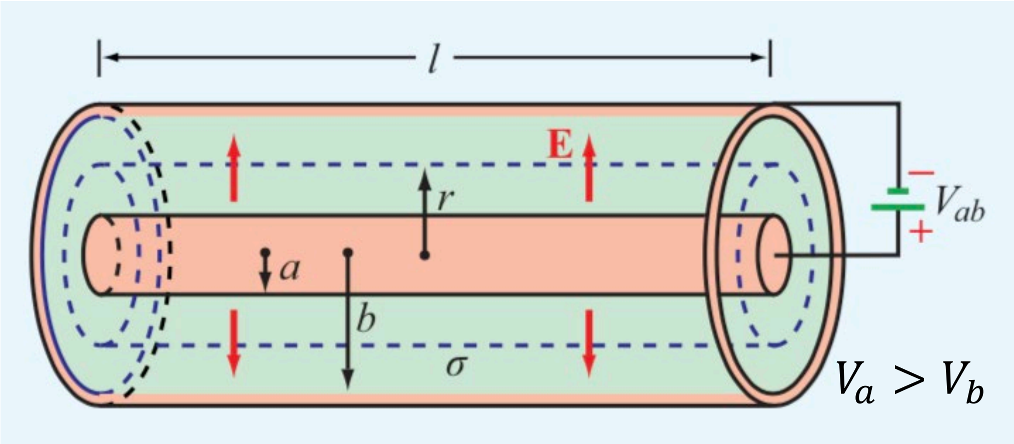
Coax cable length, l

Inner radius = a
Outer radius = b

Insulation material has conductivity, σ

Obtain G'/ unit length of insulation layer

Example - coax cable



<u>Conductance of coax</u>

Voltage inner conductor > outer

Coax cable length, l

Inner radius = a

Outer radius = b

Insulation conductivity, σ

<u>Obtain</u> G'/ unit length of <u>insulation layer</u>

I = current flowing from inner conductor to outer conductor through insulation → $\hat{r}$

Area through which current flows: $\underbrace{\mathrm{A} = 2\pi r l}_{\text{cylindrical}}$

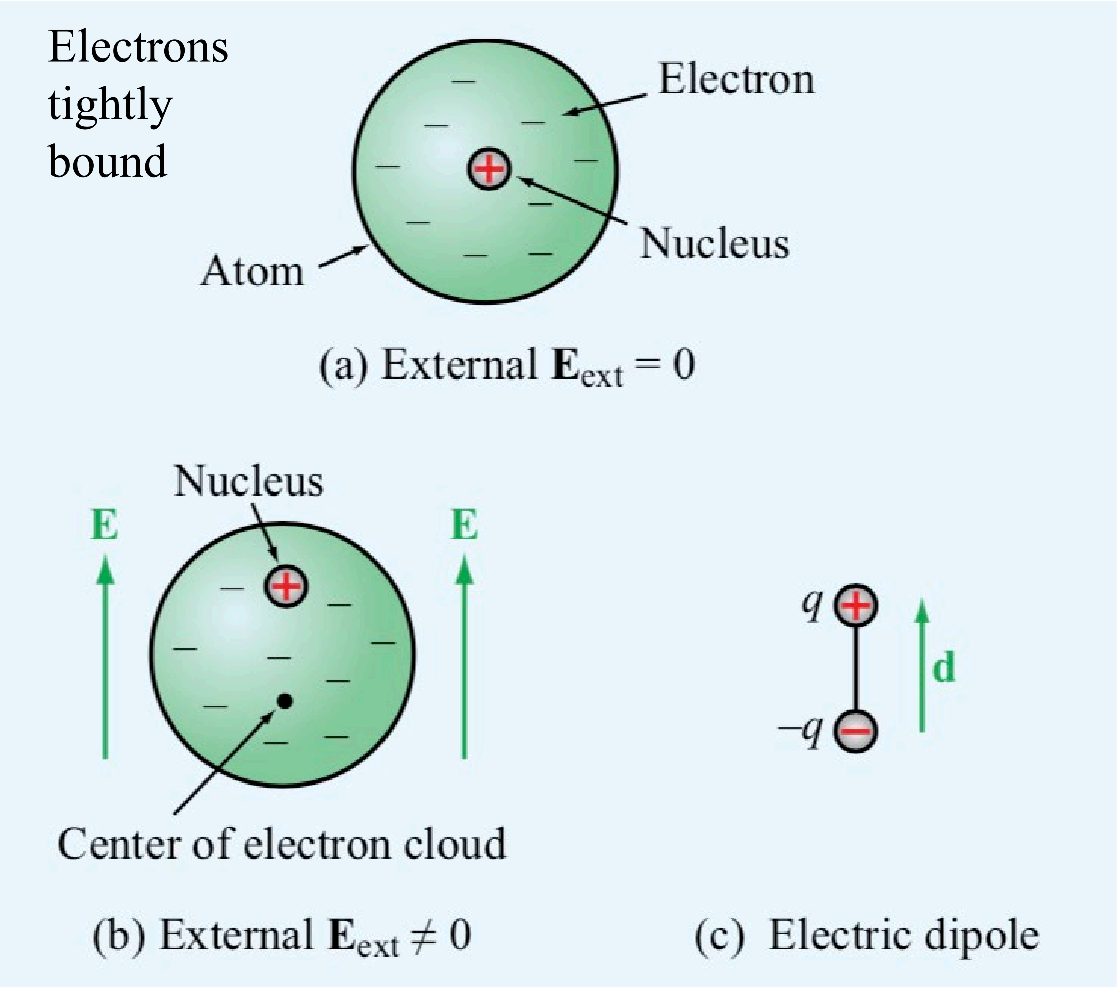
$$\vec{J} = \hat{r}\frac{I}{A} = \hat{r}\frac{I}{2\pi r l} \qquad \vec{J} = \sigma\vec{E} \qquad \vec{E} = \hat{r}\frac{I}{2\pi\sigma r l} \text{ from inner to outer}$$

$$V_{ab} = -\int_b^a \vec{E}\cdot d\vec{l} = -\int_b^a \hat{r}\frac{I}{2\pi\sigma r l}\cdot\hat{r}dr = \frac{I}{2\pi\sigma l}\ln\left(\frac{b}{a}\right)$$

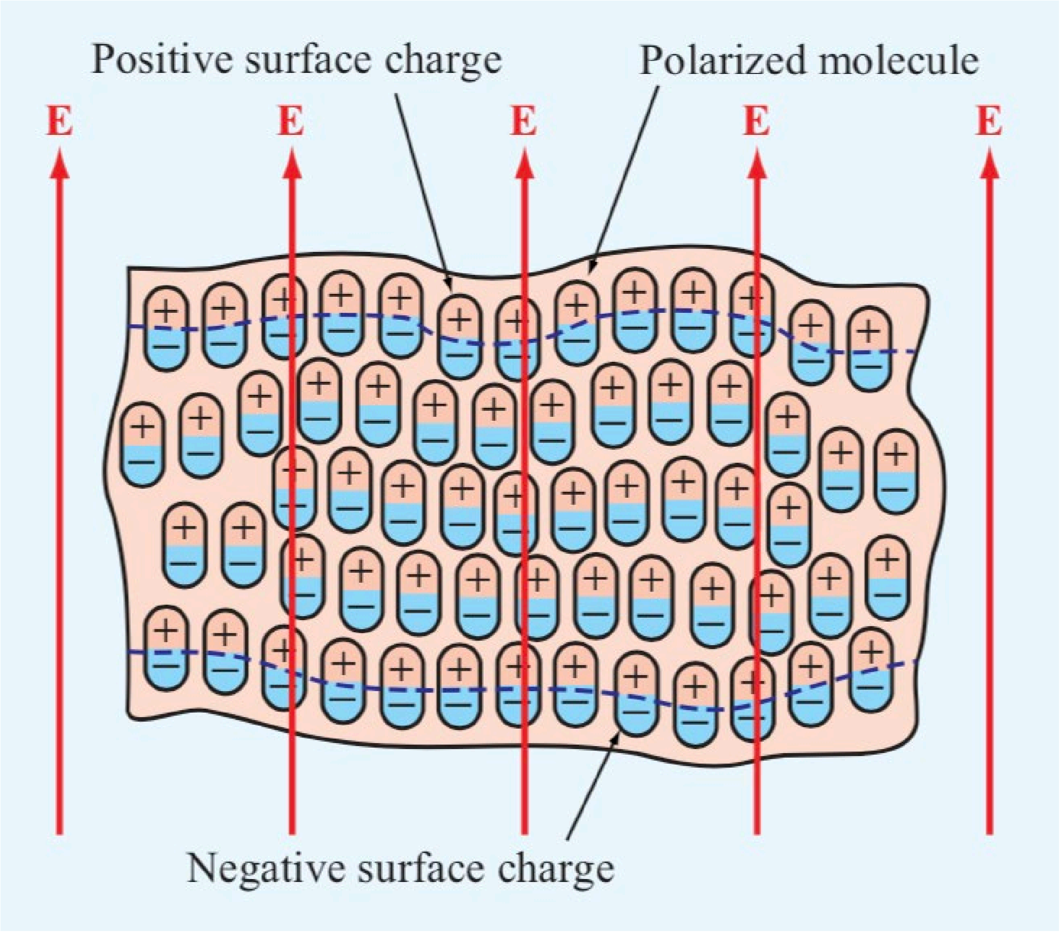
$$G' = \frac{G}{l} = \frac{1}{Rl} = \frac{I}{V_{ab}l} = \frac{2\pi\sigma}{\ln(b/a)} \; [S/m]$$

Same as table 2.1

Dielectrics – dipole model



Dielectrics



→ Can have polar materials with <u>permanent dipole moments</u> that are either random or aligned with $\vec{E}$ (liquid crystals)

Electric Breakdown

The dielectric strength E_{ds} is the largest magnitude of **E** *that the material can sustain without breakdown.*

Table 4-2: Relative permittivity (dielectric constant) and dielectric strength of common materials.

Material	Relative Permittivity, ε_r	Dielectric Strength, E_{ds} (MV/m)
Air (at sea level)	1.0006	3
Petroleum oil	2.1	12
Polystyrene	2.6	20
Glass	4.5–10	25–40
Quartz	3.8–5	30
Bakelite	5	20
Mica	5.4–6	200

$\varepsilon = \varepsilon_r \varepsilon_0$ and $\varepsilon_0 = 8.854 \times 10^{-12}$ F/m.

Polarization field

Free space $\vec{D} = \epsilon_0 \vec{E}$

In a dielectric we define: $\vec{D} = \epsilon_0 \vec{E} + \vec{P}$ ← Electric polarization field

Linear – $\vec{P}$ is linear with $\vec{E}$

Isotropic – $\vec{P}$ and $\vec{E}$ in same direction

Anisotropic – $\vec{P}$ and $\vec{E}$ may have different directions

Homogeneous - ϵ, μ, σ are uniform constant

Electric susceptibility

$$\vec{P} = \epsilon_0 \chi_e \vec{E}$$

$$\vec{D} = \epsilon_0 \vec{E} + \epsilon_0 \chi_e \vec{E} = \epsilon_0 (1 + \chi_e) \vec{E} = \epsilon \vec{E}$$

$$\epsilon = \epsilon_0 (1 + \chi_e) \qquad \epsilon_r = \frac{\epsilon}{\epsilon_0}$$

Electric boundary conditions

Gauss $\vec{\nabla} \cdot \vec{D} = \rho_V$ → $\oint_S \vec{D} \cdot d\vec{s} = Q$

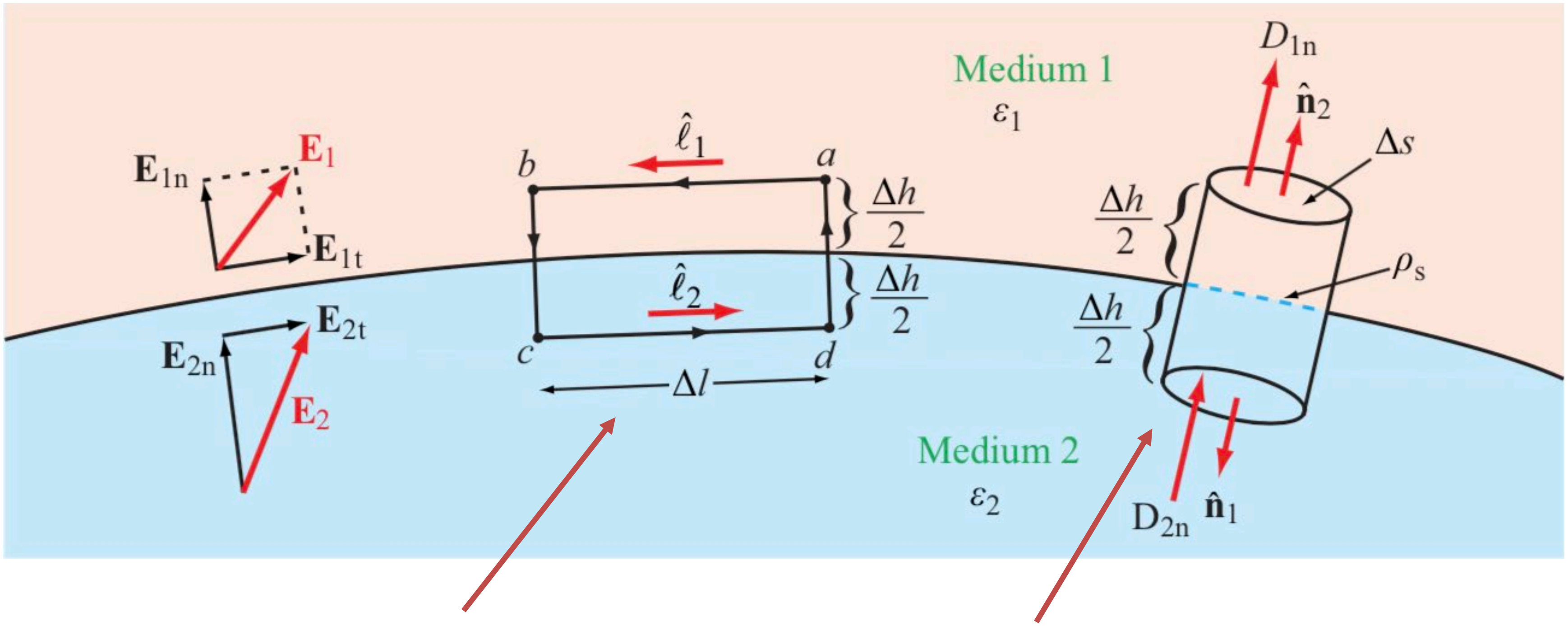
Kirchoff (Faraday): $\vec{\nabla} \times \vec{E} = 0$ → $\oint_C \vec{E} \cdot d\vec{l} = 0$

From Maxwell's equations – obtains set of boundary conditions on:
$\vec{E}$, $\vec{D}$, and $\vec{J}$ at interface of any 2 media

(later will do $\vec{H}$, $\vec{B}$)

Derived from electrostatics → apply when $\frac{\partial}{\partial t} \neq 0$

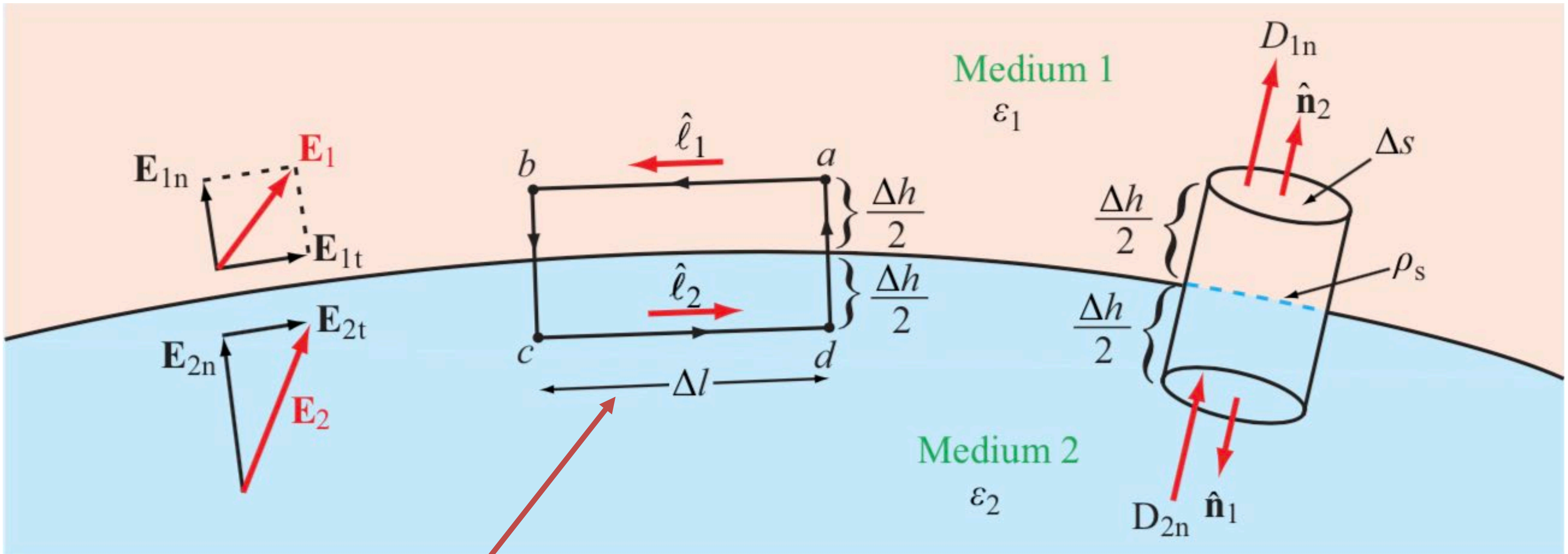
Tangential/normal components of $\vec{E}$ and $\vec{D}$



Faraday: $\vec{\nabla} \times \vec{E} = 0 \rightarrow \oint_C \vec{E} \cdot d\vec{l} = 0$

Gauss: $\vec{\nabla} \cdot \vec{D} = \rho_V \rightarrow \oint_S \vec{D} \cdot d\vec{s} = Q$

Tangential components of $\vec{E}$ and $\vec{D}$



Consider closed rectangular loop **a-b-c-d-a**

Conservative field: $\oint_C \vec{E} \cdot d\vec{l} = 0$ Line integral around closed path is zero

Let $\Delta h \to 0$ then $\overline{bc}$ and $\overline{da} \to 0$

$$\oint_C \vec{E} \cdot d\vec{l} = \int_a^b \vec{E}_1 \cdot \hat{l}_1 dl + \int_c^d \vec{E}_2 \cdot \hat{l}_2 dl = 0$$

Medium 1 Medium 2

Tangential components of $\vec{E}$ and $\vec{D}$

$$\left.\begin{array}{l} \vec{E}_1 = \vec{E}_{1t} + \vec{E}_{1n} \\ \vec{E}_2 = \vec{E}_{2t} + \vec{E}_{2n} \end{array}\right\}$$ Tangential and normal components to the boundary

$\hat{l}_1 = -\hat{l}_2 \qquad (\vec{E}_1 - \vec{E}_2) \cdot \hat{l}_1 = 0$

To satisfy Faraday's law (closed contour, conservative field):
Component of $\vec{E}_1$ along $\hat{l}_1$must equal component of $\vec{E}_2$ along $\hat{l}_1$ for all $\hat{l}_1$ tangential to the boundary

$$\boxed{\vec{E}_{1t} = \vec{E}_{2t}} \qquad \text{[V/m]}$$

Tangential components of $\vec{E}$ and $\vec{D}$

Tangential component of the $\vec{E}$ field is <u>continuous</u> across boundary for <u>any</u> 2 media

$$\vec{D}_1 = \epsilon_1 \vec{E}_{1t} + \epsilon_1 \vec{E}_{1n}$$

$$\vec{D}_2 = \epsilon_2 \vec{E}_{2t} + \epsilon_2 \vec{E}_{2n}$$

$\Rightarrow$

$$\vec{D}_{1t} = \epsilon_1 \vec{E}_{1t}$$

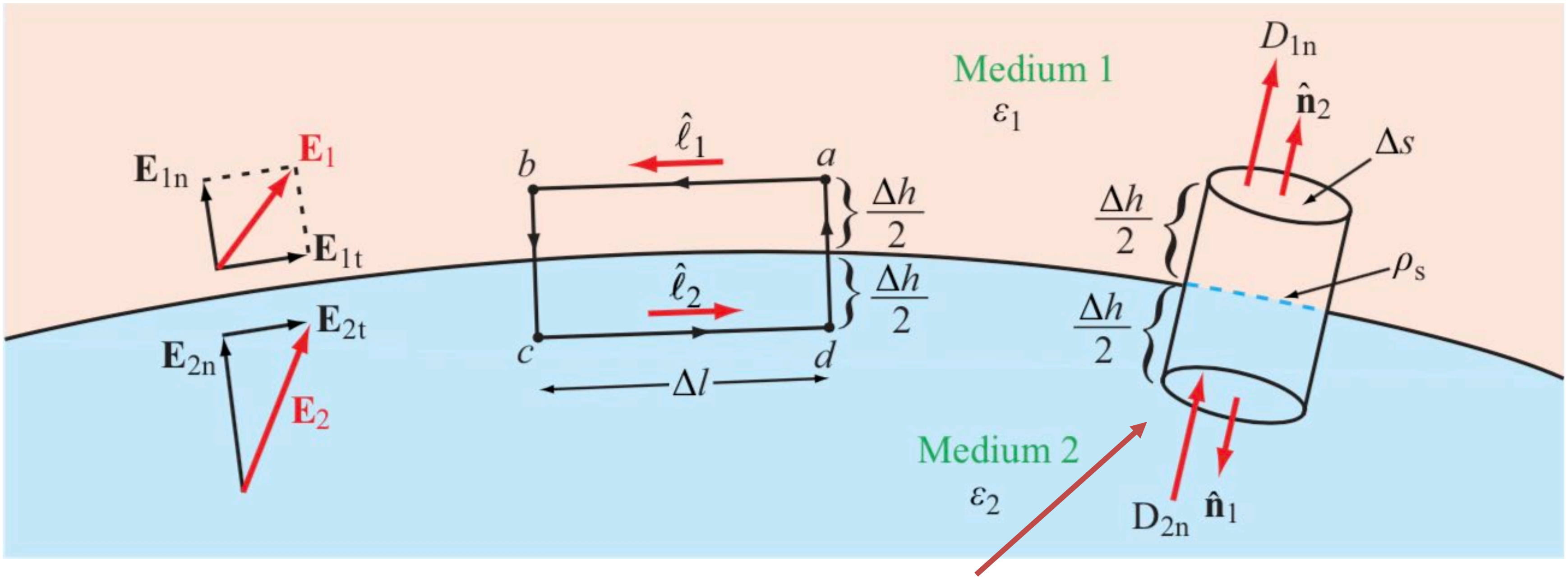
$$\vec{D}_{2t} = \epsilon_2 \vec{E}_{2t}$$

$$\vec{D}_{2t} = \epsilon_2 \vec{E}_{1t}$$

(since $\vec{E}_{1t} = \vec{E}_{2t}$)

$$\frac{\vec{D}_{1t}}{\epsilon_1} = \frac{\vec{D}_{2t}}{\epsilon_2} \qquad \frac{\vec{D}_{1t}}{\vec{D}_{2t}} = \frac{\epsilon_1}{\epsilon_2}$$

Normal components of $\vec{E}$ and $\vec{D}$



Apply Gauss's Law: total outward flux through cylinder must equal total charge inside.

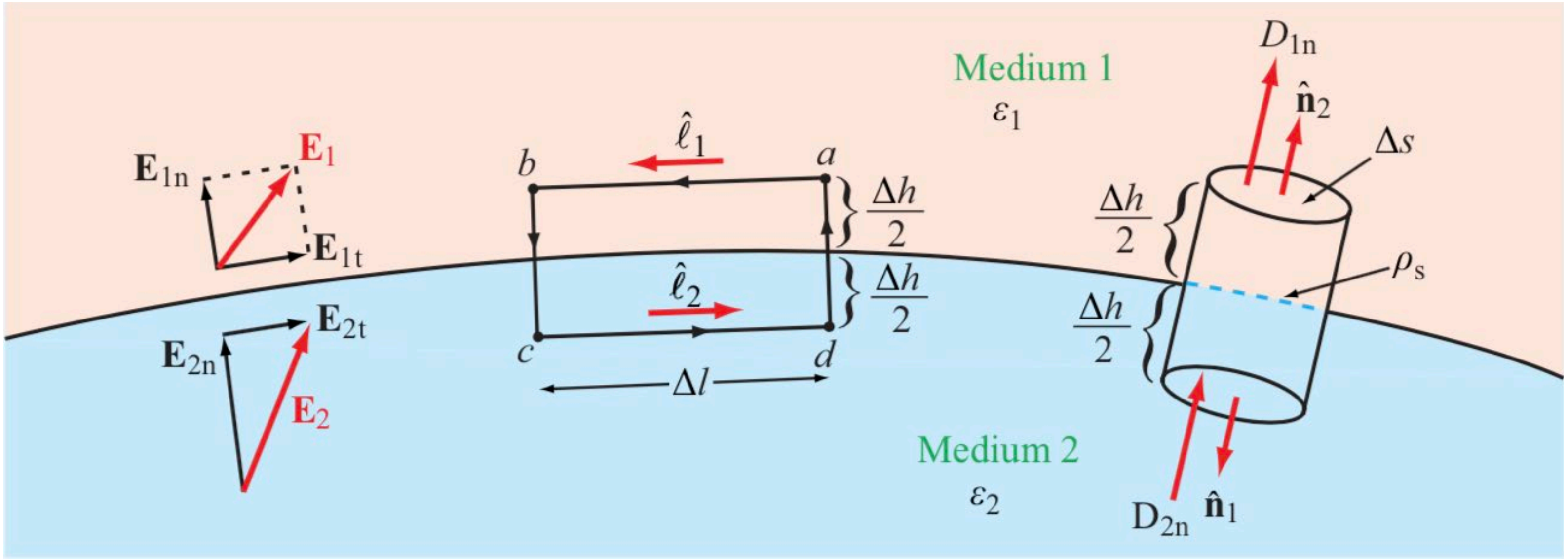
$$\oint_S \vec{D} \cdot d\vec{s} = Q \qquad \vec{\nabla} \cdot \vec{D} = \rho_V$$

Let cylinder height $\Delta h \to 0$ → only flux is from top/bottom surfaces.

$Q = \rho_s \Delta s$ Surface charge density × surface differential

$$\oint_S \vec{D} \cdot d\vec{s} = \int_{top} \vec{D}_1 \cdot \hat{n}_2 ds + \int_{bottom} \vec{D}_2 \cdot \hat{n}_1 ds = \rho_s \Delta s$$

Normal components of $\vec{E}$ and $\vec{D}$



$\hat{n}_1$ and $\hat{n}_2$ are outward normal unit vectors from the surface

$$\hat{n}_1 = -\hat{n}_2$$

$$\hat{n}_2 \cdot (\vec{D}_1 - \vec{D}_2) = \rho_s \qquad [\mathrm{C/m^2}]$$

$$\boxed{D_{1n} - D_{2n} = \rho_s} \qquad [\mathrm{C/m^2}]$$

Normal components of $\vec{D}$ change by surface charge density

$$\hat{n}_2 \cdot (\epsilon_1 \vec{E}_1 - \epsilon_2 \vec{E}_2) = \rho_s \qquad \boxed{\epsilon_1 E_{1n} - \epsilon_2 E_{2n} = \rho_s}$$

Summary

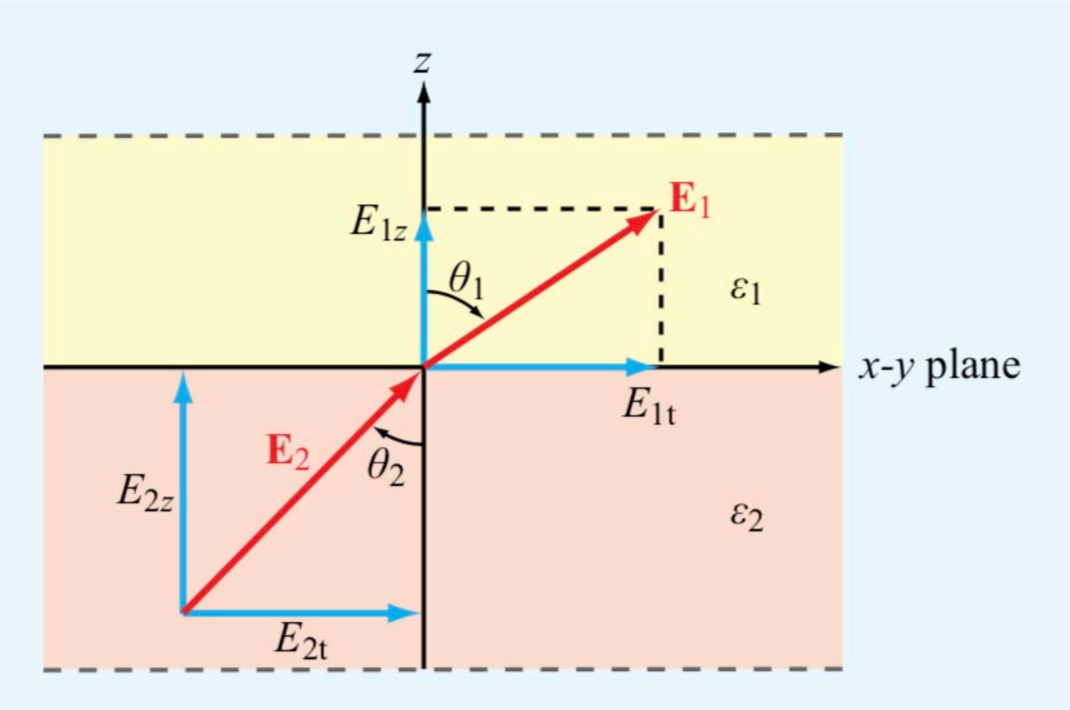
Conservative property of $\vec{E}$:

$$\vec{\nabla} \times \vec{E} = 0 \longrightarrow \oint_C \vec{E} \cdot d\vec{l} = 0 \qquad \boxed{\vec{E}_{1t} = \vec{E}_{2t}}$$

Divergence of $\vec{D}$:

$$\vec{\nabla} \cdot \vec{D} = \rho_V \longrightarrow \oint_S \vec{D} \cdot d\vec{s} = Q \qquad \boxed{\epsilon_1 E_{1n} - \epsilon_2 E_{2n} = \rho_s}$$

Example: 2 dielectrics

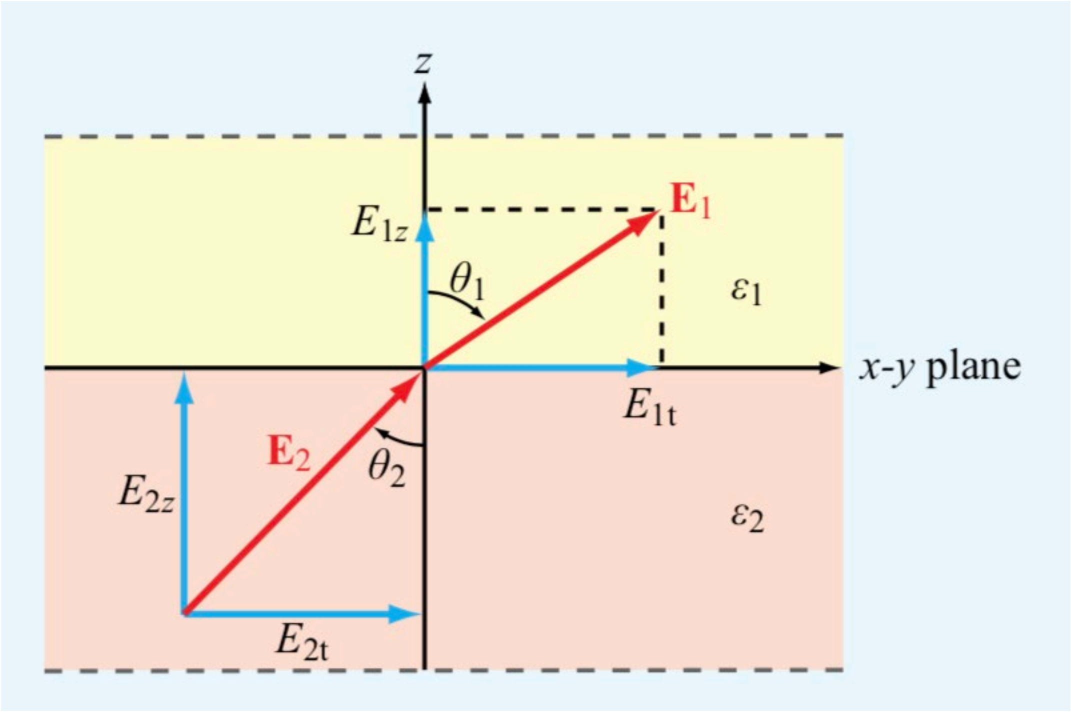


x-y plane boundary charge-free between 2 dielectrics, ϵ_1 and ϵ_2

$\vec{E}_1 = \hat{x}E_{1x} + \hat{y}E_{1y} + \hat{z}E_{1z}$
in dielectric ϵ_1

Find $\vec{E}_2$, θ_1, θ_2

Example: 2 dielectrics



x-y plane boundary charge-free between 2 dielectrics, ϵ_1 and ϵ_2

$$\vec{E}_1 = \hat{x}E_{1x} + \hat{y}E_{1y} + \hat{z}E_{1z}$$
in dielectric ϵ_1

Find $\vec{E}_2, \theta_1, \theta_2$

Let $\vec{E}_2 = \hat{x}E_{2x} + \hat{y}E_{2y} + \hat{z}E_{2z}$ Normal to interface is $\hat{z}$, x-y are tangential

$E_{2x} = E_{1x}$ and $E_{2y} = E_{1y}$ tangential

$D_{2z} = D_{1z} \rightarrow \epsilon_2 E_{2z} = \epsilon_1 E_{1z}$ normal (charge free)

$$\boxed{\vec{E}_2 = \hat{x}E_{1x} + \hat{y}E_{1y} + \hat{z}\frac{\epsilon_1}{\epsilon_2}E_{1z}}$$

Example: 2 dielectrics

To obtain angles:

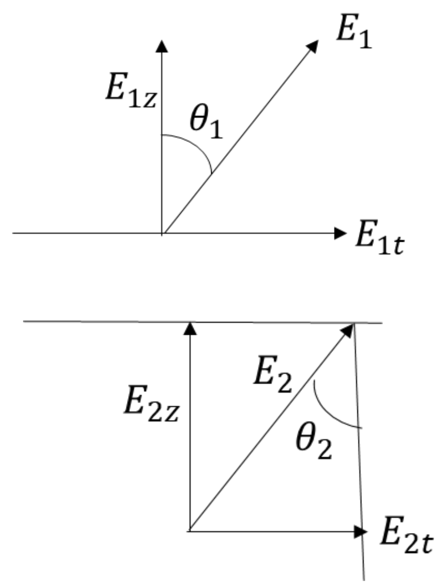
$$E_{1t} = \sqrt{{E_{1x}}^2 + {E_{1y}}^2}$$

$$E_{2t} = \sqrt{{E_{2x}}^2 + {E_{2y}}^2}$$

$$tan\theta_1 = \frac{E_{1t}}{E_{1z}} = \frac{\sqrt{{E_{1x}}^2 + {E_{1y}}^2}}{E_{1z}}$$

$$tan\theta_1 = \frac{E_{1t}}{E_{1z}} = \frac{\sqrt{{E_{1x}}^2 + {E_{1y}}^2}}{E_{1z}}$$

$$tan\theta_2 = \frac{E_{2t}}{E_{2z}} = \frac{\sqrt{{E_{2x}}^2 + {E_{2y}}^2}}{E_{2z}}$$



$$\frac{tan\theta_2}{tan\theta_1} = \frac{\epsilon_2}{\epsilon_1}$$

Dielectric – conductor boundary

Perfect conductor medium 2:

$$\vec{E}_2 = \vec{D}_2 \quad = 0$$

$$E_{1t} = D_{1t} = 0$$

$$D_{1n} = \epsilon_1 E_{1n} \quad = \rho_s \longrightarrow \boxed{\vec{D}_1 = \epsilon_1 \vec{E}_1 = \hat{n}\rho_s}$$

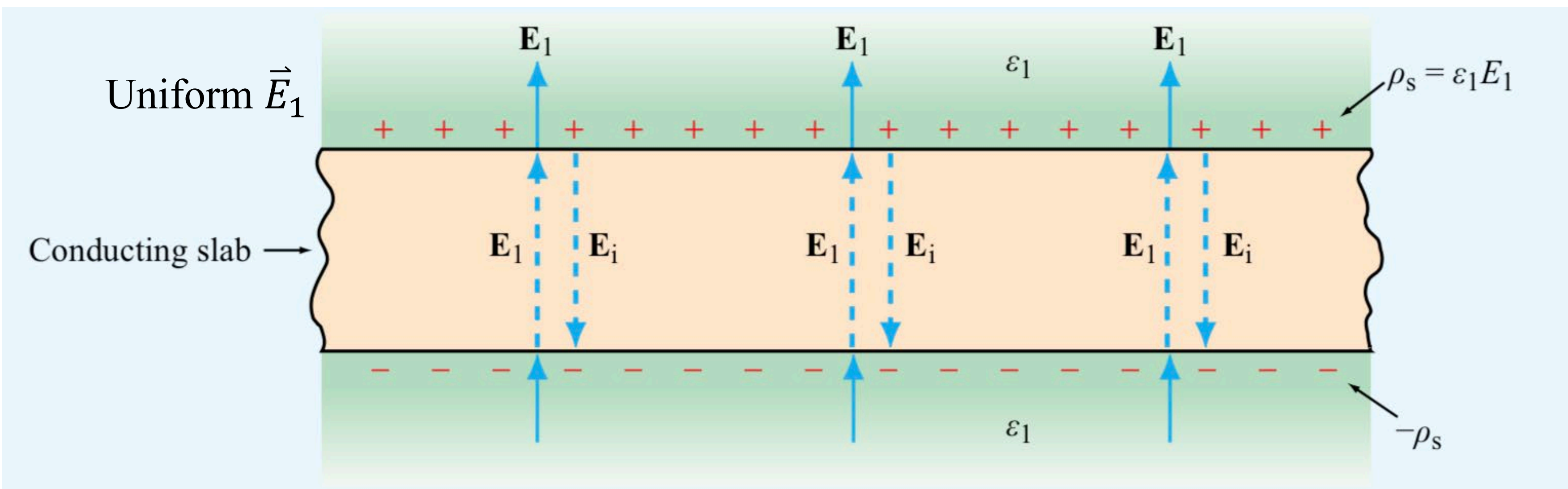
At conductor surface

$\hat{n}$ from conductor

$\hat{n}$

Medium 1: ϵ_1

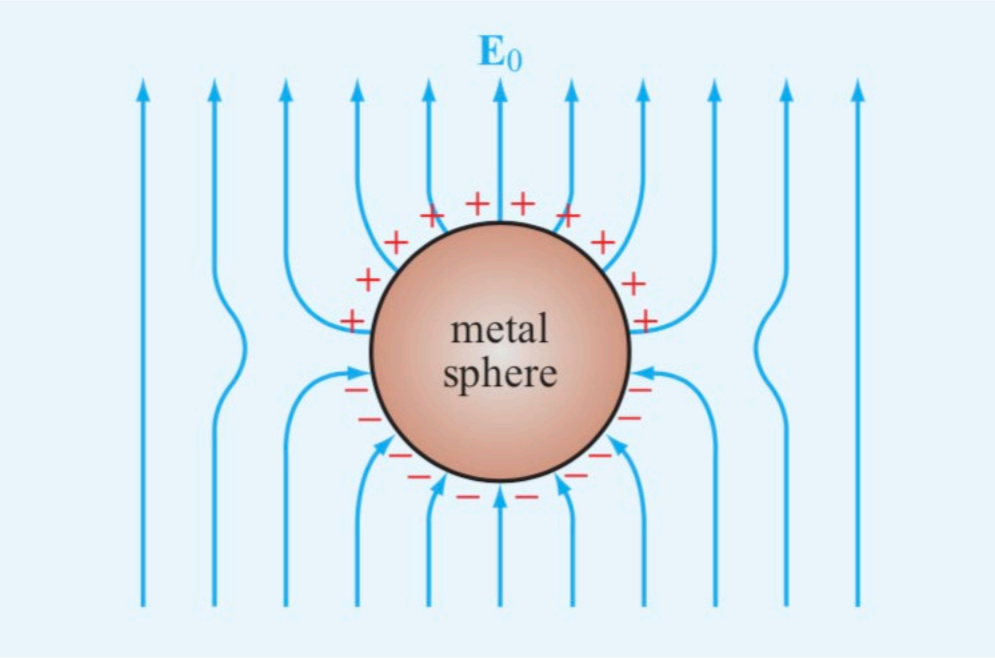
Medium 2: $\sigma = \infty$



$\vec{E}_i = -\vec{E}_1$ since field inside conductor <u>must = 0</u>

Net electric field inside a conductor is zero

Dielectric – conductor boundary



Metal in external field, $\vec{E}_0$

$\vec{E}$ points in = neg charge

$\vec{E}$ points out = pos charge

$\vec{E}$ <u>always</u> normal to conductor

Summary of Boundary Conditions

Table 4-3: Boundary conditions for the electric fields.

Field Component	Any Two Media	Medium 1 Dielectric ε_1	Medium 2 Conductor
Tangential E	$\mathbf{E}_{1t} = \mathbf{E}_{2t}$	$\mathbf{E}_{1t} = \mathbf{E}_{2t} = 0$	
Tangential D	$\mathbf{D}_{1t}/\varepsilon_1 = \mathbf{D}_{2t}/\varepsilon_2$	$\mathbf{D}_{1t} = \mathbf{D}_{2t} = 0$	
Normal E	$\varepsilon_1 E_{1n} - \varepsilon_2 E_{2n} = \rho_s$	$E_{1n} = \rho_s/\varepsilon_1$	$E_{2n} = 0$
Normal D	$D_{1n} - D_{2n} = \rho_s$	$D_{1n} = \rho_s$	$D_{2n} = 0$

Notes: (1) ρ_s is the surface charge density at the boundary; (2) normal components of $\mathbf{E}_1$, $\mathbf{D}_1$, $\mathbf{E}_2$, and $\mathbf{D}_2$ are along $\hat{\mathbf{n}}_2$, the outward normal unit vector of medium 2.